BSc and MSci EXAMINATIONS (MATHEMATICS)

May 2021

This paper is also taken for the relevant examination for the Associateship.

XXX

XXX (Solutions)

1. (a) The Vandermonde matrix and its inverse are sime sime sime seen \parallel sim. seen \Downarrow

$$
V = \begin{pmatrix} 1 & 1 & 1 & 1 \\ 0 & 1 & 0 & 1 \\ 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad V^{-1} = \begin{pmatrix} 1 & -1 & -1 & 1 \\ 0 & 1 & 0 & -1 \\ 0 & 0 & 1 & -1 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \tag{1}
$$

which can be computed by hand by e.g. doing back-substitution on the columns of the identity matrix. The basis is then

$$
\psi_1(x, y) = 1 - x - y + xy = (1 - x)(1 - y),\tag{2}
$$

$$
\psi_2(x, y) = x - xy = x(1 - y),\tag{3}
$$

$$
\psi_3(x, y) = y - xy = y(1 - x),\tag{4}
$$

$$
\psi_4(x, y) = xy.\t\t(5)
$$

(b) (i) Suitable nodal variables are $N_i(p) = p(z_i)$ where $z_1 = (1,0), z_2 = (0,1),$ $z_3 = (0,0), z_4(1/3,1/3).$

> We can check that the basis is a nodal one for these nodes by noticing that the spanning set for P is the linear functions plus a cubic "bubble" function that vanishes on the triangle vertices (and the edges). Thus to make a nodal basis for our nodal variables, the basis functions 1 to 3 can just be the usual linear basis functions (that are equal to 1 on the corresponding vertex and 0 at all the others) plus a scalar multiple of the bubble function so that they vanish at z_4 . The bubble vanishes at all the vertices so it just needs to be scaled appropriately to take the value 1 at z⁴ as required. 6, A

(ii) We assign N_i to vertex z_i for $i = 1, 2, 3$, and the bubble function the entire sim. seen \Downarrow cell. This is a C0 geometric decomposition because: (1) the value at each vertex can be obtained from the nodal variable assigned to that vertex (since it is just point evaluation at the vertex), (2) the value at each edge can be obtained from the nodal variables assigned to the closure of the edge, which is just vertex values at each end in this case, and the function is linear when restricted to an edge. $\boxed{6, B}$

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- 2. (a) The theorem is insufficient because $|u|_{H^3(K_1)}$ is unbounded, so it doesn't provide $\qquad \overline{\rm{unseen}\Downarrow}$ any bound on the error. $\boxed{6, B}$
	- (b) The appropriate statement is (under the same conditions as 5.28 but with $u \in$ $H^2(K_1)$, $k = 3$),

$$
|\mathcal{I}_{K_1}u - u|_{H^1(K_1)} \le C_1 |u|_{H^2(K_1)}.
$$
\n⁽⁶⁾

To prove it,

$$
|\mathcal{I}_{K_1}u - u|_{H^1(K_1)} \le ||Q_{3,B}u - u||_{H^1(K_1)}^2 + ||\mathcal{I}_{K_1}(u - Q_{3,B}u)||_{H^1(K_1)}^2 \tag{7}
$$

$$
\le (1 + C^2)|u|_{H^2(K_1)}^2,
$$
 (8)

where $Q_{3,B}$ is the degree k averaged Taylor polynomial over a ball B inside K_1 but as large as possible, and where we used Lemmas 3.22 and Corollary 3.16. $\vert 7, B \vert$

(c) The appropriate statement is (under the same conditions as 5.30 but with $u \in$ $H^2(\Omega)$

$$
|\mathcal{I}_K u - u|_{H^1(\Omega)} \le Ch|u|_{H^2(\Omega)}.
$$
\n
$$
(9)
$$

To show this, note that we can obtain the local estimate

$$
|\mathcal{I}_K u - u|_{H^1(K)} \le C_K d |u|_{H^{k+1}(K)},\tag{10}
$$

by following the steps in the proof but with k replaced by 1. Then the same technique of summing over all the cells gives the global result. $\boxed{7, C}$

sim. seen \Downarrow

3. (a) To derive the variational form, we multiply by a test function v and integrate by sim. seen \Downarrow parts as usual to get

$$
\int_{\Omega} \nabla v \cdot \nabla u \, \mathrm{d}x - \int_{\partial \Omega} v \underbrace{\frac{\partial u}{\partial n}}_{=g} \mathrm{d}S = 0, \tag{11}
$$

so a suitable variational form is to find $v\in \bar V_h$ such that

$$
\int_{\Omega} \nabla v \cdot \nabla u \, \mathrm{d}x = \int_{\partial \Omega} v g \, \mathrm{d}S, \quad \forall v \in \bar{V}_h,\tag{12}
$$

where $\bar V_h$ is the subspace of V_h of functions that integrate to zero, and V_h is some choice of C^0 finite element space. $\boxed{6,\, {\rm A}}$

(b) The issue is that the integrals are not tractable in general, so we can't evaluate the RHS of the problem. A possible modification is to interpolate g to V_h in the boundary resulting in g_h , and solve the perturbed problem

$$
\int_{\Omega} \nabla v \cdot \nabla u \, \mathrm{d}x = \int_{\partial \Omega} v g_h \, \mathrm{d}S, \quad \forall v \in \bar{V}_h. \tag{13}
$$

(c) The modification to Céa's Lemma is

$$
||u - u_h||_{H^1(\Omega)} \le (1 + M/\gamma) \sup_{v \in V_h} ||u - v||_{H^1(\Omega)} + \frac{C}{\gamma} ||g - g_h||_{L^2(\partial \Omega)}, \qquad (14)
$$

so there are now two terms, a best approximation term of u in V_h , and an approximation error term for g_h .

To prove it, following the steps of Céa's Lemma, we take a test function $v \in V_h$ in both the exact and approximate equation, and compute the difference, to yield

$$
a(u - u_h, v) = \int_{\partial \Omega} v(g - g_h) \, dS, \,\forall v \in V_h. \tag{15}
$$

Then we use coercivity to write (for any $v \in V_h$)

$$
\gamma \|u_h - v\|_{H^1(\Omega)} \le a(u_h - v, u_h - v),
$$
\n(16)

$$
= a(u_h - u, u_h - v) + a(u - v, u_h - v), \tag{17}
$$

$$
= \int_{\partial\Omega} (u_h - v)(g_h - g) \, \mathrm{d}S + a(u - v, u_h - v), \tag{18}
$$

$$
\leq C||u_h - v||_{H^1(\Omega)}||g_h - g||_{L^2(\partial\Omega)} + M||u - v||_{H^1(\Omega)}||u_h - v||_{H^1(\Omega)},
$$
\n(19)

where C is the constant in the trace inequality and M is the continuity constant of the bilinear form $a(u, v)$. Then, dividing by $||u_h - v||_{H^1(\Omega)}$ gives

$$
\gamma \|u_h - v\|_{H^1(\Omega)}^2 \le C \|g_h - g\|_{L^2(\partial \Omega)} + M \|u - v\|_{H^1(\Omega)}.
$$
\n(20)

Then, combining with the triangle inequality, we get

$$
||u - u_h||_{H^1(\Omega)} \le ||u - v||_{H^1(\Omega)} + ||u_h - v||_{H^1(\Omega)},
$$
\n(21)

$$
\leq (1 + M/\gamma) \|u - v\|_{H^1(\Omega)} + \frac{C}{\gamma} \|g - g_h\|_{L^2(\partial\Omega)},
$$
 (22)

and minimisation over v gives the result. $\vert 8, D \vert$

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integrating by parts separately in Ω_1 and Ω_2 gives Z $\nabla v \cdot \nabla u \, \mathrm{d} x$ $v\left(\frac{\partial u}{\partial n} |_{\partial \Omega_1} + \frac{\partial u}{\partial n} |_{\partial \Omega_2 \cap \Gamma}\right) dS = 0,$ (23)

and substitution of the boundary condition gives the variational problem: find
$$
u_h \in V_h
$$
 such that

$$
\int_{\Omega} \nabla v \cdot \nabla u_h \, \mathrm{d}x = 2 \int_{\Gamma} v \, \mathrm{d}S, \quad \forall v \in V_h. \tag{24}
$$

(b) We are in the case of Theorem 4.38, so we just need to check continuity of the linear form according to

$$
F[v] = 2 \int_{\Gamma} v \, dS \le ||v||_{L^{2}(\Gamma)} 2|\Gamma| \le ||v||_{H^{1}(\Omega_{0})} 2|\Gamma| \le ||v||_{H^{1}(\Omega)} 2|\Gamma|.
$$
 (25)

where we have used the trace theorem for continuous finite elements (Theorem 4.4), and $|\Gamma| = \int_{\Gamma} dS$. Hence F is continuous. $\boxed{8, A}$

(c) The bound studied in the course is

$$
||u_h - u||_{H^1(\Omega)} \le h|u|_{H^2(\Omega)},
$$
\n(26)

but the solution has a jump in the first derivative across Γ , so $|u|_{H^2(\Omega)}$ is not finite, so the bound does not imply convergence of the numerical solution as $h \to 0$.

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4. (a) Multiplying by a test function v that vanishes on the exterior boundary and \vert unseen \Downarrow

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5. (a) The map is surjective, so there exists $v \in V$ such that $q = \nabla \cdot v$ for all $q \in Q$. unseen \Downarrow Hence, using the Riesz Representation Theorem, for all $F\in Q'$, there exists q_F such that

$$
F[p] = \int_{\Omega} pq_F \, \mathrm{d}\, x, \quad \forall p \in Q. \tag{27}
$$

So, for all $F \in Q'$, there exists $v \in V$ such that

$$
b(v, p) = \int_{\Omega} \nabla \cdot vp \, dx = F[p], \quad \forall p \in Q.
$$
 (28)

In other words, for all $F \in Q'$ there exists v such that $Bv = F$, which means that B is surjective. Then, from the notes, this implies the inf-sup condition. $\vert 5, M \vert$

(b) Let $q \in \text{Ker}(\delta)$, i.e. $\delta q = 0$. Taking w such that $\nabla \cdot w = q$, we have

$$
0 = \int_{\Omega} \nabla \cdot wq \, dx = \int_{\Omega} q^2 \, dx \implies q = 0. \tag{29}
$$

(c) Let $q \in \text{Ker}(\delta_h)$. Then for $w \in V$,

$$
\int_{\Omega} w \cdot \delta q \, \mathrm{d} x = b(w, q) = b(\Pi_h w, q), \tag{30}
$$

$$
= \int_{\Omega} \Pi_h w \cdot \delta_h q \, \mathrm{d} \, x = 0,\tag{31}
$$

so $q \in \text{Ker}(\delta)$ as required.

(d) Considering $p \in Q_h$ having a pattern of the type given in Figure 2, it suffices to consider $b(w, p)$ for w being basis functions associated with vertices in the interior of the mesh, which span V_h (because of the zero boundary condition). The support of w consists of 6 triangles with symmetry about a diagonal line from top-left to bottom-right. The divergence of w is antisymmetric about that line, whilst p is symmetric, so the integral $b(w, p)$ vanishes. Hence, $p \in \text{ker}(\delta_h)$. The previous result says that the Fortin Trick assumptions imply that the $\ker(\delta_h)$ is empty, V_h, Q_h must fail to satisfy the Fortin Trick assumptions.

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$$

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Review of mark distribution:

Total A marks: 34 of 32 marks Total B marks: 19 of 20 marks Total C marks: 13 of 12 marks Total D marks: 14 of 16 marks Total marks: 100 of 80 marks Total Mastery marks: 20 of 20 marks